Computational Aeroelasticity: Success, Progress, Challenge

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Introduction

EARLY 70 years ago, Theodorsen¹ published a now famous report outlining, in detail, an analytical method by which the flutter characteristics of airfoils with two or three degrees of freedom could be theoretically calculated. This development laid the groundwork for what has become a rich area of research in theoretical aeroelasticity.

Recently, the term computational aeroelasticity (CAE) has been coined, and the term generally is used to refer to the coupling of highlevel computational fluid dynamics (CFD) methods to structural dynamics tools to perform aeroelastic analysis. However, given the large body of research that has led us to this point, it is inappropriate to constrain the definition of CAE in this fashion.

This article asserts a much broader definition of the term, one in which CAE encompasses all levels of aeroelastic analysis. Aeroelastic tools based on both linear unsteady aerodynamics and nonlinear CFD methods have been developed and successfully applied. We refer to both of these methodologies as components of CAE. Likewise structural modeling as simple as beam theory to state-of-the-art finite element modeling (FEM) have been incorporated into aeroelastic tools, and these techniques should also be included under the CAE heading.

It is not the intention of this article to provide an exhaustive history of the development and application of CAE over the past 70 years. Rather, the subject will be examined in the context of two primary themes: 1) aeroelastic problems requiring theoretical



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investigation and 2) strategies and techniques attacking these problems.

The discussion will be broken into three distinct categories. Problems that are relatively well understood and associated CAE tools that are viewed as mature and accepted for these types of problems will be discussed as the successes of CAE. Emerging, less understood aeroelastic problems where CAE has been applied on a limited basis, or a one-of-a-kind application of CAE, will be examined as progress in CAE. Finally, problems which continue to contest CAE methodology and roadblocks to the future development and application of CAE will be discussed as the challenges to CAE.

This categorization provides a reasonable roadmap for the future development of CAE. The successes provide us with a template for future development. They represent the tools that the user community is willing to employ on a day-to-day basis, and the problems that are important to the development of current aerospace vehicles. The progress categorization illustrates the types of problems that are beginning to limit the development and/or performance of vehicles and analysis techniques that can be developed by stretching the current technology. The challenge categorization provides a target or focus for future development. Even with today's advanced methodologies and computational capability, there are problems that cannot be accurately or efficiently modeled using current techniques, particularly in the area of nonlinear unsteady aerodynamics.

Each of these areas will be discussed in detail with specific examples illustrating the various assertions. Finally, our view of the grand challenge to the future development of CAE methodology will be outlined and discussed.

Successes in CAE

Despite the complexity of coupling three distinct engineering disciplines, aerodynamics, structures, and dynamics into a unified aeroelastic analysis capability, computational aeroelasticity has enjoyed a significant number of successes over its course of development. Today, every crewed vehicle that flies through our atmosphere undergoes some level of aeroelastic analysis before flight. Virtually every major uncrewed flight vehicle is similarly analyzed.

Flutter is a catastrophic aeroelastic phenomenon that must be avoided at all costs, and all flight vehicles must be clear of flutter and many other aeroelastic phenomena in their flight envelope. Flight and wind-tunnel testing are two ways to clear a vehicle for flutter, but both are expensive and occur late in the design process. Therefore, engineers rely heavily on computational methods to assess the aeroelastic characteristics of flight vehicles. The successes of CAE are rooted in this aeroelastic characterization process.

Problem Formulation

When examining the subject of CAE, one cannot overlook the simple elegance of the formulation of the problem. The development of the generalized aeroelastic equations of motion is enabling for most modern CAE tools and should be viewed as one of the true successes in CAE. The generalized aeroelastic equations of motion are given by

$$[M]\{\ddot{q}(t)\} + [D]\{\dot{q}(t)\} + [K]\{q(t)\} = \{F(t)\}$$
 (1)

$$\{w(x, y, z, t)\} = \sum_{i=1}^{N_{\text{modes}}} q_i(t) \{\phi_i(x, y, z)\}$$
 (2)

where $\{w(x,y,z,t)\}$ is the structural displacement at any position and time on the vehicle and $\{q(t)\}$ is the so-called generalized displacement vector, both of which are simply geometric properties describing the time history of the aeroelastic deformation. [M], [D], and [K] are the generalized mass, damping, and stiffness matrices, respectively, and ϕ_i represents the normal modes of the structure. These terms result purely from the structural and mass properties of the vehicle. The $\{F(t)\}$ term is the generalized force vector and represents the coupling of the unsteady aerodynamics and inertial loads with the structural dynamics. Thus, the coupled aeroelastic equations of motion comprise distinct terms that can be related to

the structures, aerodynamics, and dynamics disciplines that are required to formulate the problem. This allows great flexibility in the choice of methods that can be used to model a given system. For instance, for linear structural models, the generalized mass, damping, and stiffness matrices are constants that do not vary with time or the structural deformation. This remains true independent of the aerodynamics chosen for representation of the generalized force. Thus, the structural and aerodynamic models under this formulation remain completely independent of each other. Varying levels of fidelity, sometimes termed variable fidelity modeling, of structural and aerodynamic modeling can be readily matched to the problem under analysis without changing the overall formulation of the equations of motion.

This mix and match characteristic is exploited quite regularly through the choice of the unsteady aerodynamic simulation used to construct the generalized force term. This term has been modeled using aerodynamics methods that range from linear doublet lattice and kernel function solutions in the frequency domain to solution of the three-dimensional unsteady Reynolds averaged Navier–Stokes (RANS) equations in the time domain. When both the structures and aerodynamics are linear, solution of the generalized aeroelastic equations of motion reduces to computation of the complex eigenvalues of the stability matrix, whose values determine the stability of the system. Computations involving nonlinear aerodynamics and/or structures are typically performed in the time domain, which tend to complicate the process of determining system stability, but nevertheless stem from the same formulation of the problem.

This unique formulation of the equations of motion also facilitates the systematic evaluation of an aeroelastic system. Virtually all aeroelastic analyses begin with an analysis of the system stability that involves the use of linear structural and aerodynamic models. As critical points in the aircraft design and/or flight envelope are identified, the analysis can be, and regularly is, refined through incorporation of higher fidelity structural or aerodynamic models, as appropriate. This approach to the analysis provides the designers with an improved confidence that the design will be free from aeroelastic anomalies throughout the flight envelope. In addition, the approach provides designers with important data regarding the sensitivity of the design to nonlinear effects.

Subsonic/Supersonic Aeroelastic Analysis of General Aircraft Configurations

Linear subsonic and supersonic aeroelastic analysis of general aircraft configurations has matured over the last 20 years, and these types of computations are performed routinely today. Tools are now commercially available that allow aeroelastic modeling and analysis to be performed directly from existing structural FEM. Trimmed static aeroelasticity, flutter, divergence, gust response, and aeroservoelastic response are among the simulations that can be readily computed using these tools.

Linear analysestypically involve modeling complex systems with either simplified geometric or physical properties. In the case of linear aeroelastic analysis, the physics of the problem are modeled at a lower fidelity so that the problem can be linearized. Geometric approximations are usually introduced into the aeroelastic models as well, but, in general, these methods are capable of modeling relatively complete, geometrically complex configurations. The majority of modern linear aeroelastic methods are highly developed tools that provide the user with a broad range of functionality and analysis options. The methods are computationally efficient, making them expedient as rapid analysis and multidisciplinary design and optimization (MDO) tools. They also have the desirable property that classical problems with exact solutions can usually be easily modeled, significantly simplifying the verification tasks for newly developed methods.

Two types of linear aeroelastic analysis will be highlighted here. The first describes examples of linear flutter analysis using these techniques. The second discusses linear aeroservoelastic and gust response simulations.

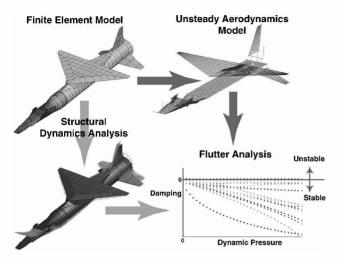


Fig. 1 X-43A flutter analysis.

Linear Flutter Analysis

An example flutter analysis of the X-43A research and launch vehicle is shown in Fig. 1. In this case the complete X-43A hypersonic research and launch vehicles were modeled structurally using FEM techniques. The model for this portion of the analysis is shown in the upper-left portion of Fig. 1. A normal modes structural dynamics analysis was performed to obtain rigid-body and structurally flexible mode shapes and frequencies as shown in the lower-left portion of Fig. 1. The model for the unsteady aerodynamics portion of the analysis is shown in the upper-right corner. For this analysis, doublet lattice² aerodynamics was used in the subsonic flight regime, whereas the ZONA51 (Ref. 3) linear unsteady aerodynamics methodology was used in the supersonic regime.

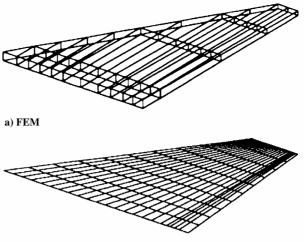
Mode shapes from the structural dynamics analysis are interpolated onto the aerodynamics grid using a surface splining technique.⁴ This allows the generalized aerodynamic force term of Eq. (1) to be computed directly from the unsteady aerodynamic analysis. Finally, the resulting flutter analysis obtained at a fixed Mach number is shown in the lower-right corner. Figure 1 outlines the general procedure used for most of today's linear flutter and aeroelastic analyses. There are a number of proprietary and off-the-shelf packages available to perform this type of analysis, and the choice of method is primarily one of user preference. Thus, linear flutter analysis is substantially viewed as a mature science.

Linear Aeroservoelastic and Gust Response Analysis

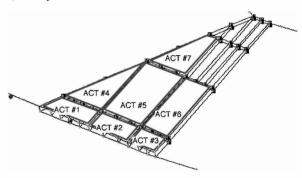
Aeroservoelastic (ASE) and gust response analyses are also performed on a relatively regular basis using linear aeroelastic analysis methods. For these applications, external forcing terms are added to the generalized force term of Eq. (1), a term describing the generalized force due to control deflection and a term describing the generalized force due to a gust. The equations are typically solved in state-space form and are constructed from separate models of the aeroelastic system, the vehicle dynamics sensors, the control actuators, the control system, and the gust spectrum. The generalized force terms are expressed as rational functions of Laplace variables in the *s* domain. This is done, for example, using the classical Rogers approximation⁵ or the minimum-state method of Karpel, ⁶ which is applied presently.

Two ASE examples are presented. The first demonstrates flutter suppression on an F-16-like wing, and the second investigates gust response reduction for a modified F/A-18 aircraft. A modified version of the Automated Structural Optimization System (ASTROS) is used as an ASE analysis tool for both examples. The method, ASTROS*, seamlessly interfaces with the ZAERO linear aerodynamic and ASE modules.

Figure 2 shows the structural, aerodynamic, and controls models for the flutter suppression analysis of this wing. The control system for this flutter suppression system consists of seven piezoelectric



b) Aerodynamic model



c) Control actuators

Fig. 2 Finite element, aerodynamic, and control actuator models of an F-16-like wing (Ref. 7).

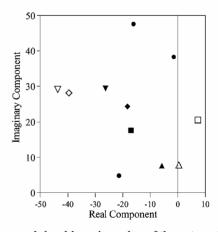
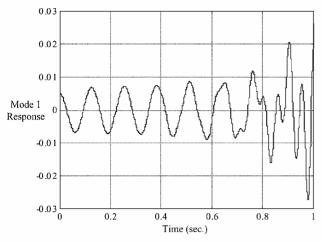


Fig. 3 Open- and closed-loop eigenvalues of the system at design air-speed, Mach = 0.9 (Ref. 7): \triangle , mode 1 (open); \square , mode 2 (open); \bigcirc , mode 3 (open); ∇ , mode 4 (open); $^{\bigcirc}$, others (open); \blacktriangle , mode 1 (closed); \blacksquare , mode 2 (closed); \blacklozenge , mode 3 (closed); \blacktriangledown , mode 4 (closed); and \spadesuit , others (closed).

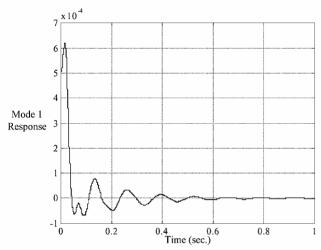
(PZT) actuators, as shown in Fig. 2. An active control system using the PZT actuators was designed to increase the open-loop flutter speed by approximately 12%.

Both open-and closed-loopflutter analyses were performed using the ASTROS* methodology and results are summarized in Fig. 3. At the design condition, the first and second modes of the open-loop system are unstable, as shown by the open square and triangle residing in the right-half plane in Fig. 3. Under the closed-loop control, these two modes move to the left-half plane as designated by the solid symbols, thus, stabilizing the system.

Figure 4 shows the modal displacement time histories for the open- and closed-loop systems after a perturbation of the system.



a) Open-loop system



b) Closed-loop system

Fig. 4 Modal time histories for the open- and closed-loop flutter suppression system (Ref. 7).

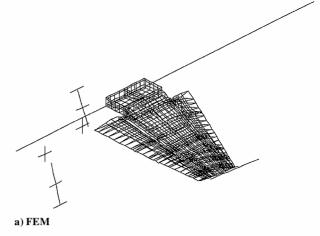
The open-loop system is clearly divergent, whereas the closed-loop system rapidly converges to a stable state.

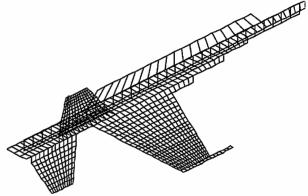
Figure 5 shows the FEM and aerodynamic models of a modified F/A-18 aircraft¹⁰ that is used to design and analyze a gust response reduction system. This model has four control surfaces: inboard/outboard leading-edge flaps, trailing-edge flap, and aileron. These control surfaces are actuated to reduce the structural response of the aircraft to an atmospheric gust.

The ASTROS* open-loopflutter analysis results in a flutter driven by the wing's torsion mode at an airspeed of 636 ft/s. The gust response reduction active control system is tailored at the design airspeed of 500 ft/s, which is 78% of the open-loop flutter speed. Figure 6 shows the rms values of the second structural mode, the torsion mode, due to a gust over a range of airspeeds. For comparison, the rms values of both open-loop and closed-loop systems are shown. The rms values of the closed-loop system are substantially reduced throughout the airspeeds of interest.

Both of these examples demonstrate that relatively complex ASE systems can be analyzed using state-of-the-art linear aeroelastic methodologies. Advanced control systems can be analyzed with and without aeroelastic effects to determine the sensitivity of the design to aeroelasticity. Gust response of aeroelstic systems can be assessed using these methods, and control systems to reduce this response can be readily evaluated. Flutter suppression control systems can be investigated, and any number of ASE phenomena can be analyzed using these techniques.

The methodology behind these analyses is among the most remarkable of the CAE success stories.





b) Aerodynamic model

Fig. 5 F/A-18 FEM and aerodynamic models for gust load alleviation analyses (Ref. 7).

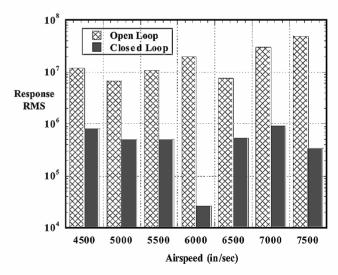


Fig. 6 Modal response rms values due to a gust vs airspeed (Ref. 7).

Transonic Wing Flutter

In the transonic speed range, aeroelastic analysis becomes significantly more complicated. Under these conditions, shock waves can form and disappear as the aircraft undergoes unsteady, structurally flexible motion. In addition, regions of separated flow can appear and disappear as these shock waves strengthen and weaken. These are highly nonlinear events that can have a profound impact on the aeroelastic behavior of flight vehicles. With flutter used as an example phenomenon, the impact of transonic unsteady aerodynamics on aeroelastic stability will be discussed.

The appearance of shock waves on the vehicle can cause a further drop in the flutter boundary in the transonic speed regime over that

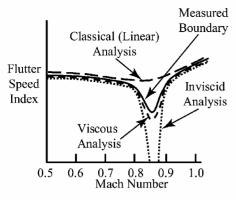


Fig. 7 Idealized flutter boundary and CAE predictions.

produced by linear compressible flow effects. This drop is termed the transonic dip. The important feature of the transonic dip is the bottom of the dip, which defines the minimum velocity at which flutter can occur across the flight envelope of the vehicle. Therefore, predicting the bottom of this dip is often crucial to the design of the vehicle.

Figure 7 presents an idealized flutter boundary through the transonic speed regime along with illustrations of typical computed results using various CAE methods. The solid line depicts the measured flutter boundary, and an expected result from a flutter analysis using linear unsteady aerodynamics is shown as the long-dashed line. As can be seen, the linear analysis typically predicts the flutter boundary adequately at subsonic and supersonic speeds, but is unconservative in the transonic speed regime, where it typically predicts a higher flutter dynamic pressure, and consequently velocity, than experiment. The flutter boundary predicted by a nonlinear, inviscidunsteady aerodynamics analysis is depicted by the dotted line. This analysis could be obtained by solving the unsteady transonic small disturbance potential flow, full potential flow, or Euler aerodynamic equations of motion. All of these methodologies have the capability of predicting shock waves in the flow, frequently resulting in a drop in the flutter boundary in the transonic speed regime. However, if viscous effects are not included in the analysis, the predicted boundary can often be overconservative, predicting a significantly lower flutter speed at the bottom of the transonic dip. Viscous effects in the form of significant boundary-layer thickening and/or shock-induced flow separation tend to define the bottom of the transonic dip, and it is only when these effects are added, as shown by the short-dashedline, that an accurate prediction of transonic flutter characteristics can be anticipated.

The severity of the dip and its Mach number range are highly dependent on the aerodynamic characteristics of the vehicle to be analyzed. Modern transports employing supercritical wing technology can experience significant transonic effects due to large unsteady motion of the shock wave across the wing chord. Wings with weaker shocks or shocks whose position is not highly sensitive to changes in local angle of attack will not experience as drastic transonic effects. These characteristics are directly related to the wing sweep, thickness, and camber distribution. Thus, the prediction of flutter characteristics in the transonic speed regime can be precarious because it is difficult, if not impossible, to predict when transonic effects will have a major impact on the flutter boundary.

The addition of nonlinearity and viscous effects to the unsteady aerodynamics analysis is not trivial. The nonlinearity of the problem typically forces the analyses to be performed in the time domain as opposed to the frequency domain for linear flows. Methods for determining the stability boundary using time-domain aerodynamics are not as simple and require a significant amount of computation and user interfacing to determine the boundary. The addition of viscous effects further complicates the process by adding longer computation times and more uncertainty in the form of turbulence modeling. Therefore, the successes in transonic flutter prediction have been limited primarily to transonic wing flutter. Modeling and com-

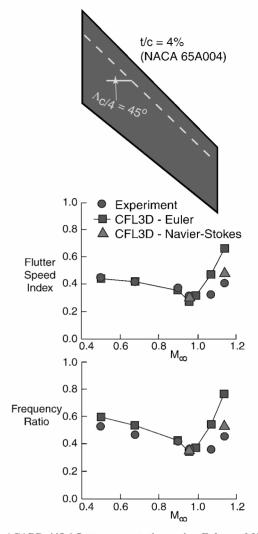


Fig. 8 AGARD 445.6 flutter computations using Euler and Navier–Stokes unsteady aerodynamics (Ref. 11).

putational requirements for more complete configuration unsteady viscous analyses have limited the widespread development and application of these types of methods to anything but the simplest of geometries. However, the advances in the use of these methods for transonic aeroelastic analysis of wings are remarkable.

Two examples of this type of application that illustrate the problems with predicting transonic aeroelasticity are presented here. The first is a series of transonic wing flutter computations by Lee-Rausch and Batina11 using Euler and Navier-Stokes unsteady aerodynamics for the AGARD 445.6 wing. The CFL3DAE code¹² was used for these computations. The planform for this wing and a summary of the flutter analysis using Euler and Navier-Stokes unsteady aerodynamics is shown in Fig. 8. The AGARD 445.6 wing has a quarter-chord sweep of 45 deg, and a symmetrical airfoil with a maximum thickness of 4%. Although there is an appreciable transonic dip for this wing, as tested in air, flutter computations using the Euler equations were able to predict the bottom of the flutter dip in the transonic speed regime, and addition of viscous effects to the analysis did not show an appreciable change in the results. The thin airfoil profile and high wing sweep combine to minimize transonic effects on this wing, resulting in this benign flutter behavior in the transonic range. At the low supersonic Mach numbers, however, the Euler equation analysis predicted a significantly higher flutter boundary, and addition of viscous effects through solution of the RANS equations improved the simulation considerably.

Gibbons¹³ demonstrated a somewhat different result in his analysis of the transonic flutter characteristics of a business jet wing using the CAP-TSD¹⁴ and CFL3DAE aeroelastic methods.

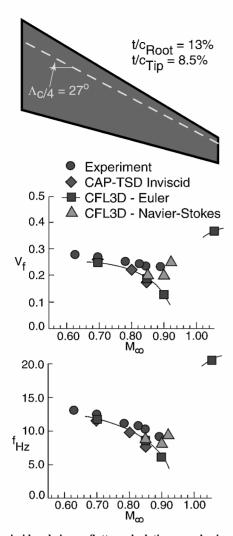


Fig. 9 Inviscid and viscous flutter calculations on a business jet wing (Ref. 13).

The wing planform and flutter results from this analysis are summarized in Fig. 9. In contrast to the AGARD 445.6 wing, this wing has a quarter-chord sweep of approximately 27 deg, and the airfoil thickness varies from 13% at the wing root to 8.5% at the wing tip. The flutter boundary shown in Fig. 9 is plotted as flutter speed index vs Mach number. Results from two inviscid analyses and a viscous analysis are compared with wind-tunnel data. For both inviscid cases, the calculated flutter boundary drops off rapidly as Mach number increases into the transonic regime. The faired line in Fig. 9 represents the CFL3D Euler computations. At Mach 0.90, the Euler computations contained enough aerodynamic damping to identify the flutter boundary shown by the square symbol. However, at this Mach number, the inviscid CAP-TSD result indicated a free response with no appreciable aerodynamic damping even for very low dynamic pressures. Euler calculations were performed at Mach 0.92 and 0.94 and showed a similar free-response characteristic. Thus, the inviscid transonic methodologies employed on this wing predict a flutter boundary that dips to very low, unrealistic velocities at transonic speeds. The addition of the viscous terms in the CFL3D computations, with appropriate changes to the wing grid, raises the predicted Mach 0.92 and 0.94 flutter boundaries to the level indicated by the triangle symbols in Fig. 9. These viscous results are in much improved agreement with the behavior observed in the wind tunnel. Similar improved agreement has been obtained with a viscous version of the CAP-TSD transonic small disturbance potential flow code.15

Although there are data that indicate that nonlinear inviscid analysis methods may be capable of predicting the transonic dip in isolated flutter cases, it would certainly seem prudent to include

viscous effects in all transonic flutter analyses. It appears that nonlinear inviscid methods produce results that are conservative in the transonic speed regime, but this has not been definitively proven. Flutter margin is often a critical parameter in aircraft design, and overly conservative flutter margins can often result in unnecessary added structural weight and/or reduced vehicle performance.

The results discussed here are enlightening from a research and phenomenological understanding standpoint and certainly demonstrate the high-quality results that can be obtained when using high-order viscous unsteady aerodynamics. However, these applications are for wing-alone geometries and have limited value in vehicle design application, where the interaction of complex geometric components can have a first-order effect on the aeroelastic performance of the vehicle. Thus, although these applications are viewed as a success, they also highlight one of the chief roadblocks to application of higher-order unsteady aerodynamics methods to aeroelastic problems. The human and computational resources required to apply these methods to problems of higher complexity than individual vehicle components precludes them from being used more extensively in general aeroelastic analysis.

Progress in CAE

The discussed successes of CAE have led to a number of isolated investigations and applications that tend to use the available CAE technology in innovative or novel ways, but have not yet received the attention to elevate the capability to mainstream application or research. This section highlights a few of these areas and discusses their role in current and future CAE development.

In general, the applications discussed here tend to stretch the available technology beyond the current accepted limits of application. In doing so, they highlight the shortcomings of the current technology and provide direction for future, more generalized development. In short, they represent the breeding ground for future CAE research and development.

Nonlinear ASE Analysis

CAE methods have been applied to a number of problems involving aeroservo and ASE analyses. CAE methods involving linear aerodynamics can be applied to some problems in this area as long as the control surface deflections are small and the geometrical discontinuities generated by activation of the control system do not generate aerodynamic nonlinearities in and of themselves. Unfortunately for most realistic control surface deflections, this is not the case, and separated flows in the vicinity of hinge lines and near control surface edges quickly degrade the accuracy of linear aerodynamic analyses. To combat these problems, researchers have applied various levels of nonlinear unsteady aerodynamic analysis methods to the prediction of unsteady airloads due to control surface deflection. Two notable examples of this type of analysis are presented here.

B-2 Residual Pitch Oscillation

During flight testing of the B-2 bomber, a nonlinear aeroelastic residual pitch oscillation (RPO) was encountered after control surface pitch doublets were input at conditions outside the aircraft's flight envelope. 16 The initial vehicle response to the control surface doublet decayed in amplitude to a small limit-cycle RPO after several pitch cycles. Video from a chase plane showed a moving shock in the condensation cloud on the upper surface of the aircraft. Further analysis of the data suggested that the RPO was a result of the interaction of the aircraft short period and first flexible bending mode of the aircraft with the oscillating shock. To better understand this RPO, an analysis effort was undertaken in which a viscous/inviscid interaction version of the CAP-TSD CAE code, CAP-TSDV,15 was modified to include the short period dynamics and the active flight control system of the aircraft. Time simulations of two conditions resulting in RPO were analyzed using CAP-TSDV with mixed results.

Figure 10 shows the B-2 planform, with the CAP-TSDV model superimposed on it. The dashed line in Fig. 10 shows the actual

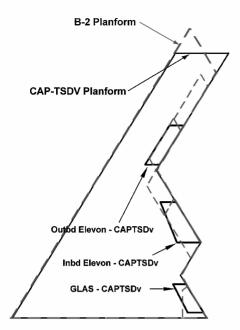


Fig. 10 CAP-TSDV model and actual B-2 planform and control surfaces (Ref. 16).

planform and control surfaces, whereas the solid line shows the CAP-TSDV representation of these components. There are a number of restrictions to the modeling capability in CAP-TSDV that precluded modeling the actual vehicle planform and control surfaces. Thus, the control surface hinge lines and edges, as well as the aircraft wing tip geometry, can only be approximated by the CAP-TSDV methodology. In this analysis, only the control surfaces actively moving during the observed RPO response were modeled, allowing the middle elevon and the outboard split drag rudders to be omitted from the model.

The first five elastic modes, as well as the rigid-body pitch and plunge modes, were modeled in the CAP-TSDV analysis. A simplified pitch control augmentation flight control system (FCS) was also included in the CAP-TSDV model to actively deflect the control surfaces during the simulation. This allowed both open- and closed-loop simulations to be performed and the impact of the control system on the RPO behavior to be assessed. Vertical acceleration and pitch rate responses computed by CAP-TSDV at the flight vehicle sensor locations were used as feedback sensor inputs for the FCS model. In addition, some of the nonlinear actuator characteristics of the flight vehicle were also modeled in the CAP-TSDV analysis.

CAP-TSDV time simulations were initiated using a converged, steady-state, trimmed analysis of the aircraft at a specified flight condition. The inboard elevon was deflected using a control surface doublet command to perturb the vehicle, and the resulting pitch response was computed by CAP-TSDV.

Both open- and closed-loop simulations were performed, and results for the closed-loop, heavy weight, forward c.g. configuration are shown in Fig. 11. The closed-loop results closely match the flight-test results at these conditions. The sharp break in the closed-loop damping characteristics has been assessed to be due to the formation of shock waves on the vehicle. The predicted frequency characteristics of the RPO, which are not shown here, also correlate with the flight-test data at these conditions.

A second, lighter weight configuration that exhibited the RPO phenomenon was also analyzed. CAP-TSDV did not predict the phenomenon as accurately as for the heavy-weight case. It overpredicted the severity of the RPO as compared to flight test and also did not capture the fact that the RPO tended to stabilize as the Mach number was further increased. Several reasons for this poorer performance have been postulated, including CAP-TSDV's ability to predict accurately the streamwise and spanwise separated flow regions at these conditions, but further research is required to formulate hard conclusions.

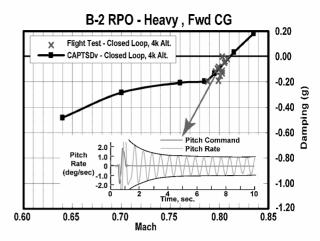


Fig. 11 Computed and flight test B-2 RPO damping characteristics (Ref. 16).

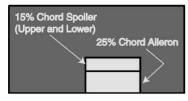
RANS Analysis of Oscillating Flap and Spoiler Configurations

The CAP-TSDV analyses of the B-2 have demonstrated the utility of viscous/inviscid interaction methods in computing fully coupled ASE problems of relatively severe complexity. However, these methods have known limitations, particularly when attempting to simulate moderately to severely separated and three-dimensional separated flows. The next step up in flow physics fidelity is to solve the RANS equations, and flow around deflected control surfaces is a prime candidate for application of this technology.

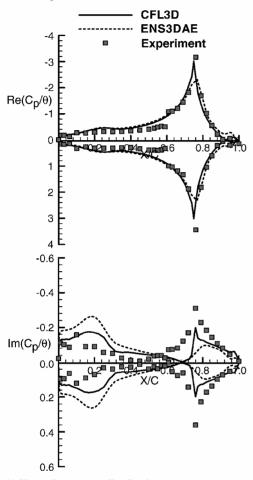
A number of studies have been performed modeling deflected control surfaces with higher-order aerodynamic methods. 17–23 References 17–20 represent a significant body of work investigating modeling and solution techniques for trailing-edge control surfaces. In these studies, particular attention has been paid to the accurate modeling of the discontinuities at the spanwise edges of control surfaces. The ENSAERO Euler/Navier–Stokes CAE method²⁴ was used for these studies, which were confined primarily to low aspect ratio, high sweep, thin wings, and wing–body configurations. The correlation with experimental data for these cases is, in general, very good, but transonic effects for these types of geometries are typically small compared to lower sweep wings with thicker wing sections.

Schuster et al.,21 Bartels and Schuster,22 and Schuster and Bartels²³ performed computations on the NASA Langley Research Center benchmark active controls technology (BACT) wing comparing two different RANS methods, CFL3D version 5.0 (Ref. 25) and ENS3DAE. 26 The BACT wing has a rectangular planform and a 12% thick symmetrical airfoil section. The planform and control surfaces for the wing are shown in Fig. 12. The two methods were compared with each other and against experimentally measured unsteady pressures along the midchord of the aileron. The analyses were performed at Mach 0.77 with the aileron oscillating sinusoidally at 5 Hz. Figure 12b shows the real and imaginary pressure coefficient plotted vs percent chord along the wing section. The computations were performed on identical grids with the only differences in the calculations being in the formulation of the equations and the turbulence modeling. The ENS3DAE calculations were performed using a central finite difference method and the Baldwin-Lomax turbulence model, whereas the CFL3D calculations were performed using an upwind finite volume method and the Spalart-Allmaras turbulence model. For the most part, the two methods give similar results with differences in the peaks at the aileron hinge line and in the imaginary component of the pressure coefficient. The imaginary component of the pressure is very small compared to the real component for this case, and differences in the imaginary component plot are overexaggerated because the vertical axis scale factor is five times that used for the real component plot.

CFL3D calculations of the wing with an oscillating spoiler have also been performed by Schuster et al.²⁷ The results of these calculations at Mach 0.77 with a mean spoiler deflection of 5 deg and a spoiler oscillation of 4.5 deg at 9.56 Hz are shown in Fig. 13.



a) BACT planform



b) Unsteady pressure distributions

Fig. 12 BACT wing planform and unsteady pressure distributions due to aileron oscillation, M = 0.77, α = 0.0 deg, δ_{ail} = 0.0 deg, θ_{ail} = 2.0 deg, and f_{ail} = 5 Hz (Ref. 22).

The computation of an oscillating spoiler geometry and quantitative comparison with experimental data is believed to be unique and represents the type of analysis that can be accurately performed using innovative modeling techniques.

Limit-Cycle Oscillation Using Viscous/Inviscid Interaction Aerodynamics

Considerable progress has been demonstrated in the use of lower-order viscous/inviscidinteraction methods to model separation onset and mildly separated flows, particularly those cases leading to limit cycle oscillation (LCO).

The LCO phenomenon can be triggered by several mechanisms. Nonlinearities in the aerodynamics, structures, or the vehicle control system can result in the quenching of an instability, or in the case of the B-2, in prevention of a convergent system from reaching a static steady state. Aerodynamic nonlinearities are difficult to model and predict due to the complex flow interactions involved. Also, the nonlinearities producing the LCO often involve a transition in flow state, such as the appearance and disappearance of shock waves, vortices, and separated flow regions.

Coupling of transonic small disturbance (TSD) potential flow methods, such as CAP-TSD, with interactive boundary-layer

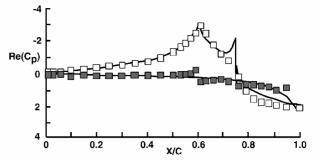
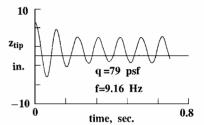
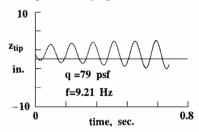


Fig. 13 Real component of unsteady pressure for the BACT wing with an oscillating upper surface spoiler, M = 0.77, $\alpha = 0.0$ deg, $\delta_s = 5.0$ deg, $\theta_s = 4.5$ deg, and f = 9.56 Hz (Ref. 27): ——, CFL3D; \square , experiment (upper surface); and \blacksquare , experiment (lower surface).



a) Amplitude decaying to LCO



b) Amplitude growing to LCO

Fig. 14 Computed business jet wing tip deflection time histories showing LCO, M = 0.89 (Ref. 28).

schemes has proven to be effective in the prediction of LCO phenomena triggered by separation onset and shock-induced separated flow. The B-2 RPO example discussed earlier is a prime example of the use of this type of methodology to predict complex LCO problems.

Edwards²⁸ has presented a second example for a business jet wing operating at transonic conditions. This case is interesting because it is purely an aerodynamics/structural LCO triggered only by nonlinear aerodynamics. The mechanism involved in this LCO is one of shock-induced separation at conditions below the transonic flutter boundary. The LCO involves the excitation of the first bending mode of the wing. Figure 14 summarizes the result of a CAP-TSDV analysis of the business jet wing at Mach 0.89 and dynamic pressure of 79 psf. In Fig. 14, the vertical deflection of the wing tip is plotted as a function of time for two different initial displacements of the wing tip. Figure 14a shows that, if the wing tip is displaced to a large deflection, the motion will decay to the experimental limit-cycle amplitude. Likewise, if the tip is initially displaced to a small deflection, the amplitude will grow until it reaches the LCO amplitude. This behavior is an important characteristic of the LCO phenomenon. The CAP-TSDV analysis captures a transient shock-induced separation on the outboard portion of the wing that is responsible for the LCO behavior. Further experiments on this wing²⁹ have verified this mechanism and identified a second transonic LCO phenomenon that involves the wing's first torsion mode.

The two LCO examples presented have demonstrated the utility of viscous/inviscid interaction methods for these types of predictions. However, the current methodology requires refinement and improvement in a number of areas.

At present, the boundary-layer formulations included in the analyses are a simple strip implementation of the two-dimensional compressible integral boundary-layer equations. The implementation of the viscous equations is also only quasi steady. Three-dimensional effects can be very important if not dominant on many modern configurations, and the two-dimensional implementation of the viscous modeling is certainly restrictive. The impact of the quasi-steady implementation of the equations has yet to be fully quantified.

On the inviscid side of the methodology, the TSD potential flow formulation of the equations limits the application of the method to flows with weak shocks and those without vortices. Higher-order inviscid formulations, such as the Euler equations, would relieve some of these restrictions.

The coupling of the methodologies, although well posed for steady flows, remains somewhat in question for unsteady flows. Addition of the viscous simulation using present techniques can have a destabilizing effect, and many times viscous computations cannot be performed, due to stability issues, in regions where similar inviscid computations can be performed.

Each of the preceding topic areas is ripe for further investigation by the research community.

Reduced-Order Modeling for Aeroelastic Applications

Finally, one cannot ignore the recent advances and contributions of reduced-order modeling (ROM) for aeroelastic and ASE applications. ROM proposes to define methodology by which physically complex dynamic phenomena can be characterized and modeled at a reduced computational cost. In addition, important physical characteristics of the system can be extracted that often cannot be identified using traditional simulation techniques. The methodology formulates strategies for identification of linear and nonlinear kernels for a given aeroelastic system. These kernels are generated using a fixed number of numerical time simulations of the system response due to a generic input, such as a pulse. These kernels can then be superimposed, through convolution, to model the system response to an arbitrary input. This type of methodology especially benefits applications such as nonlinear control system design, where the burden of many nonlinear time simulations of the complete system for each control input can be reduced to just a few simulations of the complete system for a generic input followed by many simulations using the ROM.

Beran and Silva³⁰ and Dowell and Hall³¹ provide excellent overviews of the ROM techniques under development today. A recent application by Silva and Bartels³² has demonstrated the use of RANS simulation methodology to model the transonic flutter of the AGARD 445.6 wing. Hong et al.³³ have also used this ROM technique to simulate the unsteady aeroelastic characteristics of more complex and realistic configurations. In Refs. 32 and 33, the effort required to perform a transonic flutter analysis of the wing using time simulation techniques is compared with the use of ROM.

The technique utilized in Ref. 32 requires that an aeroelastic transient due to pulse inputs in each structural mode be computed to build the ROM. In general, each of these pulse response computations are less expensive than an aeroelastic transient because the pulse responses tend to return to a steady state in fewer cycles of motion than what is required to extract frequency and damping information from a typical aeroelastic transient. Once the pulse response for each mode has been computed, the fully coupled aeroelastic response can be computed for any dynamic pressure by superimposing the impulse response characteristics through convolution. The transients computed using the ROM are relatively inexpensive to compute.

For highly complex problems where a large number of structural modes are required, the computational cost of generating modal ROM responses will increase and may surpass the cost of standard simulations. However, a primary reason for generating aeroelastic ROM models is to develop insight into the nature of the computational responses as contrasted with traditional linear aeroelastic analyses. As shown in Ref. 32, the comparison between CFD-based aeroelastic analyses can be easily performed using the CFD-based aeroelastic ROMs. In addition, methods are being developed that will enable the gen-

eration of modal-based ROMs using a single input, as opposed to an input per mode.

Thomas et al.^{34,35} present a method whereby the order reduction more directly attacks the fluid solution, thus, relieving the dependence of the ROM on the specific modal motions of the body. Their method, known as the harmonic balance technique, allows reduced-order models of the fluid problem to be developed by a superposition of a series of harmonics. These reduced-order fluid models can then be coupled to structural dynamics and control simulations to perform aeroelastic and aeroservoelastic analyses efficiently. In this case, the issue becomes how many harmonics to retain to obtain a sufficiently accurate aerodynamic analysis of the problem.

Further research and application of ROM techniques will relieve many of the issues discussed in this section, and this methodology certainly holds significant promise for future aeroelastic analysis.

Challenges to CAE

The preceding sections have discussed the areas where CAE has been used to model problems of interest to both the aircraft development and design community as well as for aeroelastic research applications. However, there are a number of areas where CAE techniques fall well short of known analysis requirements. Some of these areas will be discussed here as the challenge to CAE.

Robust, Automated Linear Flutter Analysis for Design

As an analysis tool, linear flutter methods have certainly matured to the state where they are widely applied in the aerospace industry on a day-to-day basis. Whereas great strides have been made in coupling these methodologies with mainline structural analysis tools used in the design environment, linear flutter analysis still requires a significant amount of user interaction to perform a given simulation. This lack of automation of the analysis itself has precluded its integration into the design optimization environment and it continues to be exercised as a post-design analysis tool.

Part of this problem is due to the complexity of the flutter analysis itself, but a stagnation of algorithm development focused on automating the process is probably also to blame. Output from a typical *p-k* flutter analysis for a transport wing is shown in Fig. 15. In contrast to the wing flutter analyses shown earlier, which consisted of 10 structural modes or less, this analysis contains 34 structural modes, plus 6 rigid-body degrees of freedom. Figure 15 plots modal damping vs airspeed with negative values of damping indicating aeroelastically stable motions and positive values indicating unstable motion. The critical point is when a given mode crosses the zero-dampingaxis, which defines the flutter boundary. Figure 15 is very complicated with several modes crossing the zero-damping

To incorporate flutter effectively in the design optimization process, the data of an analysis similar to that shown in Fig. 15 must be interrogated at each design condition to determine if a flutter constraint has been compromised. Simply tracking the modes in the preceding data is a formidable task, and although we have tools for this purpose, they are not infallible. For complicated systems as this,

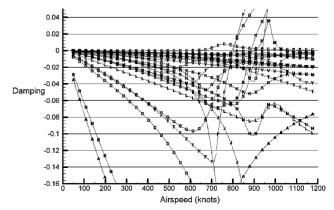


Fig. 15 Typical flutter analysis results for a transport aircraft.

it is easy to lose track of modes and inaccurately predict the critical flutter crossing.

It also is common for the flutter mechanism to change as the structural properties and even flight conditions change. Vehicles often have multiple flutter modes simultaneously active for any given flight condition. The most critical mode is the one that occurs at the lowest airspeed because the vehicle will encounterit first. However relatively small changes in the vehicle structural properties and/or flight conditions can cause flutter modes to shift dramatically, making formerly benign flutter modes critical. This can be very confusing to an automated design process because vastly different structural properties may be important to the individual flutter modes.

Further complicating the issue is that some flutter crossings are very shallow, as shown by the mode labeled E in Fig. 15. These shallow crossings can be physical flutter modes as is the case in socalled hump flutter modes, or they can be an artifice of the structural modeling. One common cause for these types of crossings is the presence of in-plane modes. For a lifting surface such as a wing, this would be a mode with deflection that is primarily in the plane of the wing, as opposed to one with deflection normal to the wing plane. Linear flutter analyses cannot accurately account for these types of motions, and the simulation can predict nonphysical flutter mechanisms. When this occurs, these in-plane modes are usually excluded from the flutter analysis to avoid the confusion caused by their presence. Aeroelasticians also often perform modal reduction to exclude modes that are not important to the flutter analysis and, thus, provide a clearer picture of the flutter characteristics of the vehicle. These types of operations are difficult to automate for use in a design environment.

Finally, the simple mathematics of the linear flutter problem can generate spurious roots to the equations that can flag false-positive results. These spurious roots are due to poor conditioning of the matrices as the problems become more complex and more modes are retained in the flutter analysis. Edwards and Wieseman³⁶ discuss many of these issues when describing the development of the Generalized Aeroelastic Analysis Method (GAAM). Automatically detecting and eliminating these roots from consideration is a must for incorporation of linear flutter methods in the automated design process.

In short, linear flutter analyses require a higher level of automation and further improvements in robust operation before they can be realistically incorporated in an MDO environment. The computational requirements of these methods are well suited to today's computer resources, and efficiency is not a major issue preventing this assimilation. At present, the entire flutter analysis process requires too much user intervention and control to be effectively incorporated in the design environment.

Efficient Nonlinear Unsteady Aerodynamic Analysis

A significant deterrent to the widespread use of high-level CFD methods in aeroelastic and ASE analysis is the computational efficiency of the methods. Significant effort has been devoted to the development and application of steady CFD, and this attention to the efficiency issues has paid off with the routine use of steady CFD, throughout the industry, as a mainstream aerodynamic analysis technique. Unfortunately, unsteady aerodynamic analysis has tended to be neglected during this development period, and we are only now beginning to seriously assess the ability of many of the steady CFD codes to handle unsteady problems effectively.

In retrospect, some of the techniques used to improve the efficiency of steady CFD methods may not be amenable to the analysis of unsteady flows. For instance, it has long been known that the low-frequency transients in a CFD simulation are a primary contributor to slow convergence rates. Significant efforts have been focused on rapidly eliminating these low-frequency transients to improve convergence to a steady state. However, many aeroelastic and ASE problems involve relatively low-frequency phenomenon, and the techniques used to develop some of today's fastest steady CFD codes may make them inappropriate for unsteady analyses due to their low-frequency damping characteristics.

A second example illustrating this point is the widespread use of domain decomposition to perform parallel computations for steady CFD analyses. When the computational grid is broken into many subdomains, or blocks, CFD computations about complex configurations involving millions of grid points can be effectively spread across many computer processors. However, extending this technique to unsteady flows results in the introduction of time lags in the unsteady analysis along each of the arbitrary subdomain boundaries. These lags force the user to use subiterations as the problem is broken into smaller blocks and spread across more processors. This has a profound impact on efficiency because the addition of just a single subiteration at best doubles the computer time required to solve the problem. Solution of the unsteady flow problem becomes dependent on the number of subdomains and the manner in which the global domain is subdivided.

New parallel processing algorithms for unsteady flows must be investigated to remove the dependence of the problem on the number of processors used to solve the equations. In other words, the solutions obtained for any given unsteady flow analysis should be identical for any number of processors used to solve the equations. This requires that implicit boundary conditions be employed at subdomain boundaries when using domain decomposition, or new parallel strategies that operate at the numerical algorithm level must be developed.

Another limiting factor is the size of time step that can be taken by a given CFD code to perform an unsteady computation. Many techniques, including implicit formulations, are limited to small time steps to maintain stability. One method that has become popular among CFD developers is the incorporation of subiteration to help converge the nonlinear aerodynamic solution between time steps. Pulliam³⁷ recommends this approach to improve the accuracy of diagonal implicit methods, and Jameson,³⁸ Melson et al.,³⁹ and Rumsey et al.⁴⁰ describe methods incorporating a dual time-stepping technique and multigrid for both accuracy and efficiency improvement of traditional factored schemes.

One must be careful, however, when assessing these methods for efficiency because each subiteration within a time step is as costly or more costly than a single time step by conventional time-stepping algorithms. This issue is highlighted by Bartels and Schuster²² using the earlier discussed oscillating aileron case as an example. The results of their experience are shown in Table 1, which shows that the CFL3D aeroelastic method using subiteration could take a time step over 25 times larger than the ENS3DAE method without subiteration. However, the CFL3D calculation required 6 subiterations per time step, and the net gain in computational work was reduced to a factor of 4.3 over the ENS3DAE solution. Comparing this net gain in efficiency to the potential loss in accuracy for taking a significantly larger time step, one can quickly see that there is a complex tradeoff between accuracy and efficiency that must be weighed when employing subiteration.

Currently, unsteady computations can be an order of magnitude or more costly than a steady flow analysis. This cost has constrained most unsteady flow applications, beyond simple research demonstrations, to be performed using structured grid formulations, which are inherently more computationally efficient than unstructured grid methods. This restriction has a direct impact on the amount of human resources required to simulate a given problem, particularly for complex configurations, and ultimately limits the acceptance of the methodology by potential users outside the CFD community. Both of these issues are discussed in subsequent sections. In short,

Table 1 Impact of subiteration on efficiency of BACT oscillating aileron calculation.²²

BACT oscillating aileron analysis	ENS3DAE	CFL3D
Time steps per cycle	3300	128
Subiterations	1	6
Computational work per cycle	3300	768
CPU time per cycle, (Min)	175	35

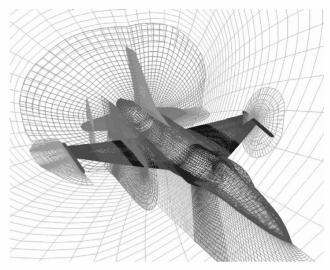
there is significant margin for further development of efficient and accurate methods for the computation of unsteady nonlinear flows. As steady method development continues to mature, it is hoped that the resources and effort devoted to steady methods will be directed toward the development of new unsteady algorithms and strategies.

Robust Moving Grid Methodology

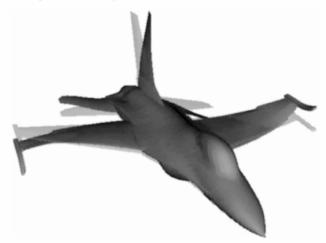
A key issue that separates the application of Euler and RANS aeroelastic methods from linear and even TSD potential flow aeroelastic schemes is the requirement of the former methods to physically move and deform the model and its surrounding grid. Linear methods eliminate the physical deformation of the vehicle from the problem by assuming small perturbation boundary conditions. TSD unsteady aerodynamics methods make a similar assumption. However, most Euler and RANS formulations employ the full surface boundary condition in their formulation, and the physical motion of the boundaries, as well as the grid surrounding the vehicle, is captured directly in the equations of motion. This is the most accurate approach to simulating the unsteady aerodynamics of the vehicle, but it introduces a major complication to the analysis. Efficiently moving what can amount to millions of grid points thousands of times during an aeroelastic calculation is a daunting task in and of itself. The fact that the grids must also be moved in such a fashion as to not degrade grid quality or introduce crossed grids or negative cell volumes adds further complication.

A significant body of research has been performed in the area of grid motion schemes for aeroelastic analysis. These range from simple algebraic shearing methods to those that model the vehicle and the surrounding grid as a physical field equation that can be coupled with the analysis. Schuster et al.²⁶ introduce a simple onedimensional algebraic shearing technique that is very efficient and is suitable for structured grid topologies. This method is effective for many geometries of interest, particularly planar lifting surface and some control surface deflection²¹ applications. However, as geometric and/or grid topology complexity increases, these simple methods begin to break down. Several researchers have extended the algebraic grid deformation approach to the use of transfinite interpolation (TFI) of the structural deformations, 41-43 which enable application to more complex block, structured grid topologies. Batina 44,45 introduced a deforming mesh algorithm based on a spring analogy that was originally developed for triangular and tetrahedral unstructured grid applications and was later extended to hexahedral cells for structured grid applications by Robinson et al. 46 Both the TFI approach and the basic spring analogy suffer from the similar character that grid skew can be exaggerated as the grid is deformed, and in extreme cases, grid crossing can occur. To combat this problem, Farhat et al.⁴⁷ recommend the use of vertex-placed torsional springs to control the angle between grid element edges. Bartels⁴⁸ further addresses these issues by imposing more strict orthogonality rules on the grid deformation in the vicinity of the vehicle surface.

Each of the described methods requires some knowledge of the grid topology and/or connectivity between grid points. Recent grid deformation research appears to be moving toward the removal of this restriction. Chen and Hill⁴⁹ describe a technique that models the volume surrounding the vehicle with an elastic homogenous solid. In this case, each grid point surrounding the vehicle is associated with a specific location in the solid, while the vehicle surface lies on the boundary of the solid. As the vehicle deforms, the elastic equations of motion are solved using a boundary element method. Each point in the grid moves according to the elastic deformation of the solid; therefore, no information concerning the structure, topology, or connectivity of the grid is required to define the grid deformation. Melville⁵⁰ describes a similar procedure that uses a distance relationship between the individual grid points and the vehicle surface to determine the grid deformation, again without knowledge of the grid topology or grid point connectivity. He has used this technique to perform an aeroelastic simulation of a complete F-16 configuration. The grid system and one of the structural modes associated with the analysis are shown in Fig. 16. As can be seen, the geometry, the associated structured, multiblock grid, and the struc-



a) Complex multiblock grid system



b) Complex structural deformation

Fig. 16 Grid and structural deformation associated with the aeroelastic analysis of an F-16 aircraft (Ref. 50).

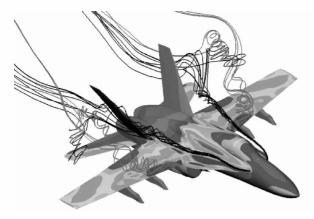
tural deformation for this analysis are all very complex. The fact that an aeroelastic simulation can be performed on a model of this complexity is a testament to the grid motion scheme employed in the analysis.

Although significant progress has been realized in the development of grid deformation schemes, no single technique has been identified as having significant benefits over the others. It appears the preprocessing approaches to the problem are gaining favor because they allow the user to assess the suitability of the grid, ahead of time, for worst-case deformations. However, there is still significant opportunity to make advances in this area of research.

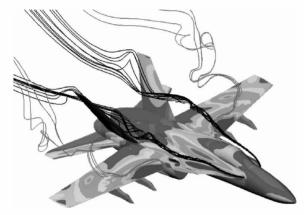
Complex Configuration Aeroelasticity

Many linear and nonlinear aeroelastic phenomena are the result of subtle structural and aerodynamic interactions between aircraft components. Engine/nacelle/pylon, wing store/pylon, strakes, chines, vortex generators, and fences are among the long list of components that can have a dramatic impact on both the steady and unsteady character of the vehicle aerodynamics and structural dynamics. It is often imperative to model these details accurately and efficiently when performing some aeroelastic analyses.

The importance of geometric details to the aerodynamic and aeroelastic performance of air vehicles is shown in Fig. 17. Sheta and Huttsell⁵¹ have numerically investigated the buffeting of the F/A-18 vertical tails due to the leading-edge extension (LEX) vortices. The effects of adding LEX fences to the aircraft are



a) LEX fences off



b) LEX fences on

Fig. 17 Effect of LEX fences on the vortex flow patterns on an F/A-18 aircraft.

clearly observed in Fig. 17. In this 35-deg angle-of-attack calculation, the LEX vortex pattern is significantly altered through addition of the fences. This pattern has a direct impact on the dynamic loads experienced on the vertical tails of this aircraft. Thus, this relatively small geometric detail has a major impact on the aeroelastic performance of the F/A-18 and other twin-vertical-tail fighters.

A second example is the LCO experienced on the F-16.⁵² This nonlinear aeroelastic phenomenon manifests itself only for certain underwing store configurations and flight conditions. It was not discovered before flight testing of the aircraft. To date, this phenomenon continues to be the subject of significant debate throughout the aeroelastic community, and potential causes ranging from complex, unsteady transonic flow interactions to structural nonlinearities have been offered as the root cause for the LCO. Unfortunately we do not possess high-order aeroelastic tools capable of accurately simulating the unsteady transonic flow about the highly complex store geometries of the F-16 to help answer some of these questions.

As these types of details are added to the simulation, the issues of the preceding two sections become enabling to the aeroelastic analysis, particularly for high-order methods. However, just the modeling requirements associated with these configuration details can be significant even for steady flow problems. For this reason, the use of unstructured grids has become popular for many steady computational aerodynamics simulations on complex configurations. The computational resources required to perform unstructured grid computations are generally higher than those required for structured grid, but the time and human resources required to model complex configurations using unstructured grids is much smaller than for structured grids.

Unfortunately, the extension of unstructured grid technology to unsteady flows has not been thoroughly investigated, and further research in this area is required. Batina⁵³ and Rausch and Batina⁵⁴

have applied an unstructured grid Euler analysis method to aeroelastic problems on complex configurations, and this work is most noted for the development of the spring analogy for grid deformation and their investigations into adaptive grid techniques for unsteady flows. Farhat et al.⁵⁵ use a similar method to perform aeroelastic analysis on an F-16 configuration. Viscous computations using unstructured grids are becoming more commonplace for steady flows, but unfortunately their extension to unsteady flows and aeroelastic problems has not yet been realized.

Nonexpert User Environment

Perhaps the most daunting challenge faced by developers of future aeroelastic analysis and design methods is the formulation of tools that can be employed by engineers who are not experts in the individual disciplines that form the methods. Aeroelasticity has a significant impact on the performance of modern air vehicles, and incorporating aeroelastic considerations early in the design process can help developers avoid design surprises in the latter stages of the aircraft development, where fixes to the problem can be very costly both financially and to vehicle performance. If these methods are integrated into the preliminary design process, engineers that are experts in fields other than CFD, finite element structural modeling, flight controls, or even aeroelasticity will use them. This will require that the methods be significantly more robust and automated than current methods, particularly for the higher-order formulations.

The developers of higher-order methods must aggressively address this situation if these tools are to be used in a fashion even similar to present linear methods. Aeroelastic analysis of flight vehicles is not optional, and organizations will take the path of least resistance to meet their aeroelastic analysis needs. At present, highorder aeroelastic methods require a significant investment in time, human, and computer resources. Until these investments are reduced, high-order aeroelastic methods will continue to be viewed as optional, and their application will be intermittent at best. Even though the cost of aeroelastic wind-tunnel and flight testing is very high, maintaining a specialized staff simply to employ high-order aeroelastic methods is higher. Aircraft manufactures will continue to maintain the status quo in regard to aeroelastic analysis until these methods become more cost effective.

Conclusions

CAE continues to play a critical role in the development of modern air vehicles. The current suite of linear aeroelastic, ASE, and gust response methodologies serves as the mainstay for aeroelastic analysis and they form a solid base on which to build future developments. Unfortunately, modern aircraft systems continually uncover aeroelastic issues that cannot be effectively predicted by these methods, and as a result, costly aeroelastic wind-tunnel and flight tests must be conducted to investigate and rectify these problems. This situation is envisioned to worsen as new concepts, such as morphing, continuous moldline controls, etc., call for vehicles with higher degrees of structural flexibility.

Many of today's analyses investigate only the aeroelastic stability of the vehicle, where small aeroelastic perturbations can be assumed and the impact of wing thickness and camber and even static aeroelastic deformation can be neglected. These characteristics will not be negligible on many future vehicles, and heretofore accepted and/or calibrated errors from linear aeroelastic analysis will no longer be acceptable.

Therefore, it is imperative that higher-order aeroelastic methods continue to be developed and refined. Many of today's problems already require the inclusion of transonic and separated flow effects, detailed structural modeling, and provisions for including structural nonlinearities. Future applications will likely require that temperature effects and chemical reactions be included in the simulations. It is our opinion that methods should continue to be developed on three levels of complexity: 1) linear methods 2) moderate fidelity methods (including viscous/inviscid interaction techniques, ROM, etc.) and 3) high-fidelity methods (RANS aerodynamics, structured and unstructured grid methods, nonlinear FEM,

etc.). During this development, consideration must be given to how the methods can be effectively employed in a design environment, not just as a postdesign analysis tool. They should also be developed with the specific objective that they be readily integrated into existing aeroelastic analysis frameworks. Finally, and most important, the new methods must be sufficiently general in application, robust and efficient in operation, and simple in implementation to make them a viable tool for general use in the aerospace industry.

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